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correction term

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## Abstract

We improve the Berezin-Li-Yau inequality in dimension two by adding a positive correction term to its right-hand side. It is also shown that the asymptotical behaviour of the correction term is almost optimal. This improves a previous result by Melas, [10].

## 1 Introduction

Let  $\Omega$  be an open bounded set in  $\mathbb{R}^d$  and let  $-\Delta$  be the Dirichlet Laplacian on  $\Omega$ . We denote by  $\lambda_j$  the non-decreasing sequence of eigenvalues of  $-\Delta$ . The main object of our interest in this paper is the lower bound

$$\sum_{j=1}^k \lambda_j \geq \frac{dC_d}{d+2} V^{-\frac{2}{d}} k^{\frac{d+2}{d}}, \quad C_d = (2\pi)^2 \omega_d^{-2/d}, \quad (1)$$

where  $V$  stands for the volume of  $\Omega$  and  $\omega_d$  denotes the volume of the unit ball in  $\mathbb{R}^d$ . Inequality (1) was proved in [8], and is commonly known as the Li-Yau inequality. In [7] it was pointed out that (1) is in fact the Legendre transformation of an earlier result by Berezin, see [1]. Note also that the Li-Yau inequality yields an individual lower bound on  $\lambda_k$  in the form

$$\lambda_k \geq \frac{dC_d}{d+2} V^{-\frac{2}{d}} k^{\frac{2}{d}}. \quad (2)$$

For further estimates on  $\lambda_k$  see [12, 5, 6, 7].

It is important to compare the lower bound (1) with the asymptotical behaviour of the sum on the left-hand side, which reads as follows:

$$\sum_{j=1}^k \lambda_j = \frac{dC_d}{d+2} V^{-\frac{2}{d}} k^{\frac{d+2}{d}} + \tilde{C}_d \frac{|\partial\Omega|}{V^{1+\frac{1}{d}}} k^{1+\frac{1}{d}} + o\left(k^{1+\frac{1}{d}}\right) \quad \text{as } k \rightarrow \infty \quad (3)$$

with

$$\tilde{C}_d = \frac{\sqrt{\pi} \Gamma\left(2 + \frac{d}{2}\right)^{1+\frac{1}{d}}}{(d+1) \Gamma\left(\frac{3}{2} + \frac{d}{2}\right) \Gamma(2)^{\frac{1}{d}}}.$$

The first term in the asymptotics (3) is due to Weyl, see [15]. The second term in (3) was established, under suitable conditions on  $\Omega$ , in [3, 4, 11], see also [13, Chap. 1.6].

It follows from (3) that the constant in (1) cannot be improved. On the other hand, since the second asymptotical term is positive, it is natural to ask whether one might improve (1) by adding an additional positive term of lower order in  $k$  to the right-hand side. The first step towards this goal was done by Melas, [10], who showed that the inequality

$$\sum_{j=1}^k \lambda_j \geq \frac{dC_d}{d+2} V^{-\frac{2}{d}} k^{\frac{d+2}{d}} + M_d \frac{V}{I} k, \quad I = \min_{a \in \mathbb{R}^2} \int_{\Omega} |x - a|^2 dx \quad (4)$$

holds true with a factor  $M_d$  which depends only on the dimension. Note however, that the additional term in the Melas bound does not have the order in  $k$  predicted by the second term in (3). Moreover, the coefficient of the second term in (3) reflects explicitly the effect of the boundary of  $\Omega$ , whereas such a dependence is not seen in the coefficient  $V/I$  of (4).

Our aim is to improve (1) and (4) by adding a positive contribution which reflects the nature of the second term in the asymptotic (3). Recently, one of the authors, see [14], proved an analogous improved estimate on the quantity

$$\sum_k (\Lambda - \lambda_k)_+^{\sigma}, \quad \sigma \geq 3/2$$

with a remainder term which agrees, up to a constant, with the corresponding second term in the asymptotics of  $\sum_k (\Lambda - \lambda_k)_+^\sigma$  as  $\Lambda \rightarrow \infty$ . The proof given in [14] relies on sharp Lieb-Thirring inequalities for operator valued potentials and works only for  $\sigma \geq 3/2$ . Since the estimates treated in present paper concern the value  $\sigma = 1$ , the method of [14] cannot be carried over to this case. We will therefore develop a different approach. We wish to point out [2] (and references therein), where an improvement of (1) has been established for the discrete Laplacian. In particular, the authors make use of a more detailed analysis of the approximation of characteristic functions by low energy eigenvalues of the discrete Laplace operator. This is also the key element of this paper.

The main idea of our strategy is explained in section 2.3. It is closely related to a modified proof of inequality (1), which we briefly describe in section 2.1, see also [9, Chap. 12]. The main results which represent improved Li-Yau inequalities in case  $d = 2$  are formulated in section 3. Since our proof includes many technical results concerning the geometry of the boundary of  $\Omega$ , we will first give its exposition for polygons, section 4. Finally, in section 5 we extend the proof to general domains.

To keep the presentation as short and stringed as possible, we have decided to restrict ourselves to the case  $d = 2$  throughout the paper.

## 2 Preliminaries

Following notation will be employed in the text. By  $\Theta(\cdot) : \mathbb{R} \rightarrow \mathbb{R}$  we denote the Heaviside function defined by  $\Theta(x) = 0$  if  $x \leq 0$  and  $\Theta(x) = 1$  if  $x > 0$ . For given  $t > 0$  we denote by  $N_t$  the number of eigenvalues of the Dirichlet-Laplacian in  $\Omega$  less than or equal to  $t$ . Finally, we will write  $[s]$  for the integer part of a real number  $s$ .

### 2.1 Li-Yau bound revisited

Let  $\psi_j$  be the sequence of the normalised eigenfunctions of  $-\Delta$  in  $\Omega$ , i.e.

$$-\Delta \psi_j = \lambda_j \psi_j \quad \text{in } \Omega, \quad \psi_j = 0 \quad \text{on } \partial\Omega, \quad \int_{\Omega} |\psi_j|^2 = 1. \quad (5)$$

In order to explain the idea which will lead to an improvement of the results by Li-Yau and Melas, it is illustrative to see how to obtain inequalities (1) and (4) for  $d = 2$  (the same arguments apply to higher dimensions as well). Following [1, 10] we extend the eigenfunctions  $\psi_j$  continuously by zero to the whole of  $\mathbb{R}^2$  so that they remain in  $H^1(\mathbb{R}^2)$ . Next introduce the following functions:

$$f_j(\xi) = (2\pi)^{-1} \int_{\Omega} e^{-ix \cdot \xi} \psi_j(x) dx, \quad F(\xi) := \sum_{j=1}^k |f_j(\xi)|^2. \quad (6)$$

Since  $\{\psi_j\}$  is an orthonormal basis of  $L^2(\Omega)$ , the Parseval identity implies that

$$F(\xi) = \sum_{j=1}^k |f_j(\xi)|^2 \leq \sum_{j=1}^{\infty} |f_j(\xi)|^2 = (2\pi)^{-2} \int_{\Omega} |e^{-ix \cdot \xi}|^2 dx = (2\pi)^{-2} V \quad (7)$$

holds for any  $\xi \in \mathbb{R}^2$ . Next we denote by  $F^*(|\xi|)$  the decreasing radial rearrangement of  $F$ . Using the well-known properties of the radial rearrangement we find

$$\int_{\mathbb{R}^2} F^*(|\xi|) d\xi = \int_{\mathbb{R}^2} F(\xi) d\xi = k \quad (8)$$

and

$$\sum_{j=1}^k \lambda_j = \int_{\mathbb{R}^2} |\xi|^2 F(\xi) d\xi \geq \int_{\mathbb{R}^2} |\xi|^2 F^*(|\xi|) d\xi. \quad (9)$$

To find a lower bound on  $\sum_{j=1}^k \lambda_j$  it thus suffices to find the minimiser of the functional  $\int_{\mathbb{R}^2} |\xi|^2 F^*(|\xi|) d\xi$  under the conditions (7) and (8).

The result of Li and Yau can be proved using the fact, [9, Chap. 12], that this functional is minimised by the function

$$\Phi_{LY}(|\xi|) = \begin{cases} (2\pi)^{-2} V & 0 \leq |\xi| \leq r_k, \\ 0 & r_k < |\xi|, \end{cases} \quad (10)$$

where  $r_k$  is given by the condition

$$(2\pi)^{-1} V \int_0^{r_k} |\xi| d|\xi| = k \Rightarrow r_k = \sqrt{\frac{4\pi k}{V}}.$$

Inserting (10) into (9) we obtain inequality (1) for  $d = 2$ .

## 2.2 Melas' improvement revisited

Melas observed in [10] that the lower bound on the right-hand side of (9) can be improved, if one takes into account that the following additional regularity condition on  $F^*$  must hold

$$|(F^*)'| \leq 2(2\pi)^{-2} \sqrt{VI} =: L. \quad (11)$$

It can be easily verified that, depending on the value of  $k$ , the corresponding minimiser  $\Phi_M$  of the functional (9) then has the following form:

$$\text{for } k \geq \frac{V^2}{48\pi I} \quad \Phi_M(|\xi|) = \begin{cases} (2\pi)^{-2} V & 0 \leq |\xi| \leq s_k, \\ (2\pi)^{-2} V - (|\xi| - s_k) L & s_k < |\xi| \leq t_k, \\ 0 & t_k < |\xi|, \end{cases} \quad (12)$$

where the points  $s_k$  and  $t_k$  are uniquely determined by

$$2\pi \int_{\mathbb{R}_+} \Phi_M(|\xi|) |\xi| d|\xi| = k, \quad t_k = s_k + \frac{V}{4\pi^2 L},$$

see Figure 1, and

$$\text{for } k < \frac{V^2}{48\pi I} \quad \Phi_M(|\xi|) = \left( \left( \frac{3kL^2}{\pi} \right)^{1/3} - L|\xi| \right)_+. \quad (13)$$

Using this minimiser we obtain the lower bound

$$\sum_{j=1}^k \lambda_j \geq \frac{2\pi}{V} k^2 + \frac{1}{32} \frac{V}{I} k \quad \text{if } k \geq \frac{V^2}{48\pi I} \quad (14)$$

and

$$\sum_{j=1}^k \lambda_j \geq \frac{2\pi}{V} k^2 + \left( 1 - 10 \cdot 2^{-\frac{5}{3}} 3^{-\frac{4}{3}} \right) \frac{3}{10} \left( \frac{2}{\pi} \right)^{\frac{2}{3}} L^{-\frac{2}{3}} k^{\frac{5}{3}} \quad \text{if } k < \frac{V^2}{48\pi I}. \quad (15)$$

Now let  $a \in \mathbb{R}^2$  be such that  $I = \int_{\Omega} |x - a|^2 dx$  and let  $B_a$  be the disc centred in  $a$  and with the volume  $V$ . It is then straightforward to verify that

$$I \geq I(B_a) = \frac{V^2}{2\pi}.$$

Using this inequality and the fact that  $k \geq 1$  we deduce from (14) and (15) the uniform estimate

$$\sum_{j=1}^k \lambda_j \geq \frac{2\pi}{V} k^2 + \frac{1}{32} \frac{V}{I} k \quad \forall k \in \mathbb{N}. \quad (16)$$

### 2.3 The new correction term

Our main observation is that the crucial reservoir for improvements of (1) does not lie in the regularity of  $F^*$ , but in a more detailed analysis and improvement of the condition (7). Indeed, since

$$F(\xi) = \sum_{j=1}^k |f_j(\xi)|^2 = \frac{V}{4\pi^2} - \sum_{j=k+1}^{\infty} |f_j(\xi)|^2, \quad (17)$$

any estimate from below on  $\sum_{j=k+1}^{\infty} |f_j(\xi)|^2$  will automatically lead to a sharper upper bound on  $F$  and therefore to an additional term in the Li-Yau inequality.

Moreover, the last term in (17) cannot go to zero arbitrarily fast as  $k$  goes to infinity. This follows from the fact that  $|e^{-ix \cdot \xi}| = 1$  everywhere in  $\Omega$ , which means that the Fourier coefficients  $f_j(\xi)$  of  $e^{-ix \cdot \xi}$  with respect to the basis  $\{\psi_j\}$  cannot decay too fast in  $j$  (each  $\psi_j$  vanishes on  $\partial\Omega$ ). In particular, the sequence  $\{f_j(\xi)\}$  is not in  $\ell^1$ . Another way to see this is to realize that the Fourier series  $\sum_j f_j(\xi) \psi_j(\cdot)$  of continuous functions approximates, in  $L^2(\mathbb{R}^2)$ , the function  $e^{-ix \cdot \xi} \chi_{\Omega}$ , which has a discontinuity on  $\partial\Omega$ . Thus the decay properties of  $\sum_{j=k+1}^{\infty} |f_j(\xi)|^2$  and consequently the additional term in Li-Yau inequality should reflect the effect of the boundary of  $\Omega$ .

The main technical difficulty is to quantify this strategy into a uniform lower bound on  $\sum_{j=k+1}^{\infty} |f_j(\xi)|^2$ . In particular, if we can prove an estimate of the form

$$\sum_{j=k+1}^{\infty} |f_j(\xi)|^2 \geq \varepsilon k^{\delta} \quad \forall \xi \in \mathbb{R}^2, \quad (18)$$

where  $\varepsilon$  and  $\delta$  are positive, then the corresponding minimiser of (9) satisfying conditions (8) and (17) reads

$$\Phi(|\xi|) = \begin{cases} V/4\pi^2 - \varepsilon k^{-\delta} & 0 \leq |\xi| \leq \tau_k, \\ 0 & \tau_k < |\xi|, \end{cases} \quad (19)$$

see Figure 1. Here  $\tau_k$  is defined by the condition

$$2\pi \int_{\mathbb{R}_+} \Phi(|\xi|) |\xi| d|\xi| = k.$$

A direct calculation then shows that there exists a positive coefficient  $A(\varepsilon, \delta)$  such that

$$\sum_{j=1}^k \lambda_j \geq 2\pi \int_{\mathbb{R}_+} \Phi(|\xi|) |\xi|^3 d|\xi| = \frac{2\pi}{V} k^2 + A(\varepsilon, \delta) k^{2-\delta}. \quad (20)$$

The asymptotic formula (3) implies that  $\delta \geq 1/2$ . For  $\delta < 1$  we obtain an improvement of the Melas bound.

Let us finally mention that a similar effect of the boundary on the sum of the eigenvalues in the case of the discrete Laplace operator was already observed in [2].

## 3 Main results

We will state and prove the results for the case of polygons and general domains separately.

### 3.1 Case 1: Polygons

For a given polygon  $\Omega$  we denote by  $p_j$ ,  $j = 1, \dots, n$  the  $j$ -th side of  $\Omega$ . Moreover, we denote by  $d_j$  the distance between the middle third of  $p_j$  to  $\partial\Omega \setminus p_j$ . We can now formulate our first result.

**Theorem 1** (Lower bound for polygons). *Let  $\Omega$  be a polygon with  $n$  sides. Let  $l_j$  be the length of the  $j$ -th side of  $\Omega$ . Then for any  $k \in \mathbb{N}$  and any  $\alpha \in [0, 1]$  we have*

$$\sum_{j=1}^k \lambda_j \geq \frac{2\pi}{V} k^2 + 4\alpha c_3 k^{\frac{3}{2}-\varepsilon(k)} V^{-\frac{3}{2}} \sum_{j=1}^n l_j \Theta \left( k - \frac{9V}{2\pi d_j^2} \right) + (1-\alpha) \frac{V}{32I} k, \quad (21)$$

where

$$\varepsilon(k) = \frac{2}{\sqrt{\log_2(2\pi k/c_1)}} \quad (22)$$

and

$$c_1 = \sqrt{\frac{3\pi}{14}} 10^{-11}, \quad c_3 = \frac{2^{-3}}{9\sqrt{2}36} (2\pi)^{\frac{5}{4}} c_1^{1/4}. \quad (23)$$

### 3.2 Case 2: General domains

For general open domains  $\Omega \subset \mathbb{R}^2$  we will have to impose certain assumptions on the regularity of  $\partial\Omega$ .

**Assumption A.** There exist  $C^2$ -smooth parts  $\Gamma_j \subset \partial\Omega$  at the boundary of  $\Omega$ . Let  $j = 1, \dots, m$ .

To be able to state the result for general domains we need some definitions. Let  $A_j, B_j$  be the end points of  $\Gamma_j$  and let  $\{x_1^j(s), x_2^j(s)\}$  be the parametrisation of  $\Gamma_j$  with its length  $s$ . We define

$$\varkappa_j = \max_s |\varkappa_j(s)|,$$

where  $\varkappa_j(s)$  denotes the curvature at the point  $s \in \Gamma_j$ . Moreover, let  $L(\Gamma_j)$  be length of  $\Gamma_j$ . Now we divide  $\Gamma_j$  into several pieces of the same length. The tiling of  $\Gamma_j$  will be done in two different ways depending on the values of  $\varkappa_j$  and  $L(\Gamma_j)$ :

(i) If

$$L(\Gamma_j) \leq \frac{3\pi}{8\varkappa_j}, \quad (24)$$

then we divide  $\Gamma_j$  into three parts of the same length and denote by  $d_j$  the distance of the middle part to  $\partial\Omega_j \setminus \Gamma_j$ .

(ii) If

$$L(\Gamma_j) > \frac{3\pi}{8\varkappa_j}, \quad (25)$$

then we divide  $\Gamma_j$  into  $n_j = \lceil 8L(\Gamma_j)\varkappa_j/\pi \rceil$  parts of the same length. Let  $a_i^j, a_{i+1}^j$  be the end points of the  $i$ -th part with  $a_0^j = A_j, a_{n_j}^j = B_j$  and let

$$\delta_i^j = \text{dist} \left( (a_i^j, a_{i+1}^j), \partial\Omega \setminus \{(a_{i-1}^j, a_i^j) \cup (a_i^j, a_{i+1}^j) \cup (a_{i+1}^j, a_{i+2}^j)\} \right)$$

Then we define

$$d_j = \min_{1 \leq i \leq n-2} \delta_i^j.$$

Finally, we will need

$$k_j := \frac{V}{2\pi} \max \left\{ \Lambda_3(j), \frac{9}{d_j^2}, \frac{128\varkappa_j^2}{\pi^2}, \frac{6\varkappa_j}{d_j} \right\},$$

where

$$\Lambda_3(j) := \max \left\{ 9 \cdot 2^{10} \max_j \varkappa_j^2, 2^{26} c_1 V^{-1}, c_1^{-1} 2^{22} 6^8 \varkappa_j^4 V \right\}.$$

Now we are in position to state the result for general domains.

**Theorem 2** (Lower bound for general domains). *Let  $\Omega$  satisfy Assumption A. Then for any  $k \in \mathbb{N}$  and any  $\alpha \in [0, 1]$  we have*

$$\sum_{j=1}^k \lambda_j \geq \frac{2\pi}{V} k^2 + \alpha c_3 k^{\frac{3}{2}-\varepsilon(k)} V^{-3/2} \sum_{j=1}^m L(\Gamma_j) \Theta(k - k_j) + (1 - \alpha) \frac{V}{32I} k. \quad (26)$$

### 3.3 Remarks

**Remark 1.** Note that the coefficient of the second term on the right hand side of (26) is very similar to the coefficient of the second term in the Weyl asymptotics (3). In particular, it reflects the expected effect of the boundary of  $\Omega$ . On the other hand, this boundary term becomes visible only for  $k$  large enough. However, we would like to point out that the second term cannot be simply proportional to  $\sum_j L(\Gamma_j)$ . Indeed, one can make  $\sum_j L(\Gamma_j)$  arbitrarily large by “folding” the boundary  $\partial\Omega$  while keeping the eigenvalues  $\lambda_j$  with  $j \leq k$  almost unchanged. This shows that the condition  $k \geq k_j$  cannot be removed.

**Remark 2.** It would be natural to try to deduce the result for general domains from the result for polygons by approximating  $\Omega$  by polygons. However, the contribution of the second term would in general disappear in such a procedure. To see this it suffices to take an open ball in  $\mathbb{R}^2$  as  $\Omega$ . Then the coefficients  $k_j$  would go to infinity when approximating  $\Omega$  by a sequence of polygons. Therefore a different strategy will be needed in the proof of Theorem 2.

**Remark 3.** As for the constants in (26), notice that  $\varepsilon(k) \ll 1$  for all  $k$  and that  $\varepsilon(k) \rightarrow 0$  as  $k \rightarrow \infty$ . On the other hand, the values of  $k_j$  are in general very large. Nevertheless, the correction term on the right-hand side of (26) can be optimised according to the geometry of  $\Omega$  by choosing the boundary segments  $\Gamma_j$  in an appropriate way.

## 4 Proof for polygons

The proofs of our main results rely on a careful exploitation of the ideas described in section 2.3.

Let  $\lambda = \lambda_k$  and let  $\mathcal{L}_k := \left\{ \sum_{i=1}^k c_i \psi_i : \sum_{i=1}^k |c_i|^2 \leq V \right\}$ . Since  $e^{i\xi \cdot x}$  belongs to  $L^2(\Omega)$  for each  $\xi \in \mathbb{R}^2$ , it follows that

$$\inf_{\psi \in \mathcal{L}_k} \|e^{i\xi \cdot x} - \psi\|_{L^2(\Omega)}^2 \leq \|e^{i\xi \cdot x} - \sum_{i=1}^k (e^{i\xi \cdot x}, \psi_i)_{L^2(\Omega)} \psi_i\|_{L^2(\Omega)}^2 = V - 4\pi^2 F(\xi), \quad (27)$$

where

$$\sum_{i=1}^k \left| (e^{i\xi \cdot x}, \psi_i)_{L^2(\Omega)} \right|^2 = 4\pi^2 F(\xi) \leq V.$$

Equation (27) yields the estimate

$$F(\xi) \leq (4\pi^2)^{-1} \left( V - \inf_{\psi \in \mathcal{L}_k} \|e^{i\xi \cdot x} - \psi\|_{L^2(\Omega)}^2 \right).$$

In view of the arguments given in section 2.3, to prove (21) it thus suffices to show that

$$\|e^{i\xi \cdot x} - \psi\|_{L^2(\Omega)}^2 \geq \text{const } k^{-\frac{1}{2}-\varepsilon(k)} \quad \forall \psi \in \mathcal{L}_k \quad (28)$$

holds for  $k$  large enough. Moreover, it is well known that  $\lambda_k \sim k$  in dimension  $d = 2$ , which shows that (28) is equivalent to

$$\|e^{i\xi \cdot x} - \psi\|_{L^2(\Omega)}^2 \geq \text{const } \lambda_k^{-\frac{1}{2}-\varepsilon(k)} \quad \forall \psi \in \mathcal{L}_k. \quad (29)$$

The idea how to prove (29) is obvious; since  $|e^{i\xi \cdot x}| = 1$  everywhere and  $\psi = 0$  on  $\partial\Omega$ , we will estimate the left-hand side of (28) by integrating over a suitable neighbourhood of  $\partial\Omega$  only. More precisely, we will make use of the contributions from integrating  $|e^{i\xi \cdot x} - \psi|^2$  over squares of the size of order  $\lambda^{-1/2}$  attached to the boundary of  $\Omega$ , see Figure 2. To estimate these contributions from below, we will need appropriate integral upper bounds on the normal derivatives of  $\psi$  on  $\partial\Omega$  in terms of  $\lambda$ . This will be done as the first step of the proof.

#### 4.1 Eigenfunctions estimates

In this section we give an  $L^2$  estimate on the derivatives the eigenfunctions  $\psi_i$  in the vicinity of  $\partial\Omega$ . Let

$$\omega = \left[0, \frac{1}{2\sqrt{\lambda}}\right] \times \left[-\frac{1}{4\sqrt{\lambda}}, \frac{1}{4\sqrt{\lambda}}\right]$$

and assume that  $\lambda$  is large enough so that the square  $\omega$  can be placed inside  $\Omega$  in such a way that one of its sides coincides with a part of  $\partial\Omega$ . We also introduce a local system of coordinates  $(x_1, x_2)$ . Finally, for a given  $p \in \mathbb{N}$  we define the sequence  $A_n(p)$  by

$$A_n(p) = (3 + 726 \cdot 4^6 p^4) A_{n-2}(p) + 150 \cdot 9^2 p^2 A_{n-1}(p) \quad (30)$$

where  $A_0(p) = 1$  and  $A_1(p) = 1$ . We then have

**Lemma 1.** *Let  $\psi_i$  be a normalised eigenfunction of the Dirichlet Laplacian on  $\Omega$  with an eigenvalue  $\lambda_i \leq \lambda$ . Then*

$$\left\| \frac{\partial^{p+1} \psi_i}{\partial x_1^{p+1}} \right\|_{L^2(\omega)}^2 \leq A_p(p) \lambda^{p+1} \quad (31)$$

holds true for all  $p \in \mathbb{N}_0$ .

*Proof.* For  $n, p \in \mathbb{N}$  we define the functions  $g : [0, 1] \rightarrow [0, 1]$  by  $g(x) := 1 - 6x^4 + 8x^6 - 3x^8$  and  $v_{n,p} : \mathbb{R} \rightarrow \mathbb{R}$  by

$$v_{n,p}(t) = \begin{cases} 1 & 0 \leq t \leq \frac{2p-n}{2p}, \\ g(2pt - 2p + n) & \frac{2p-n}{2p} \leq t \leq \frac{2p-n+1}{2p}, \\ 0 & \frac{2p-n+1}{2p} < t \end{cases}$$

with  $v_{n,p}(t) = v_{n,p}(-t)$  for  $t < 0$ . It is easy to check that

$$|v_{n,p}(t)| \leq 1, \quad |v'_{n,p}(t)| \leq 2\alpha_1 p, \quad |v''_{n,p}(t)| \leq 4\alpha_2 p^2,$$

where  $\alpha \leq 5/2$  and  $\alpha_2 \leq 11$ . Next we define

$$W_{n,p,\lambda}(x_1, x_2) = v_{n,p}(\sqrt{\lambda} x_1) v_{n,p}(4\sqrt{\lambda} x_2)$$

and note that

$$\begin{aligned} |W_{n,p,\lambda}(x_1, x_2)| &\leq 1, & |\nabla W_{n,p,\lambda}(x_1, x_2)| &\leq 9\sqrt{\lambda} \alpha_1 p \\ |\Delta W_{n,p,\lambda}(x_1, x_2)| &\leq \sqrt{2} 4^3 \lambda \alpha_2 p^2 \end{aligned} \quad (32)$$

for all  $(x_1, x_2) \in \omega$ . We will prove

$$\begin{aligned} \left\| \frac{\partial^n \psi_i}{\partial x_1^n} \right\|_{L^2(\text{supp } W_{n-1,p,\lambda})}^2 &\leq A_{n-1}(p) \lambda^n, \\ \left\| \frac{\partial^n \psi_i}{\partial x_1^{n-1} \partial x_2} \right\|_{L^2(\text{supp } W_{n-1,p,\lambda})}^2 &\leq A_{n-1}(p) \lambda^n \end{aligned} \quad (33)$$

by induction in  $n$  for  $n = 1, \dots, p$ . Notice that, in view of (58), (59), the inclusion

$$\omega_n := (\text{supp } W_{n,p,\lambda}) \subset \Omega$$

holds true for every  $p \in \mathbb{N}$  and every  $n \leq p$ . For  $n = 1$  we have

$$\left\| \frac{\partial \psi_i}{\partial x_1} \right\|_{L^2(\omega_0)}^2 \leq A_0(p) \lambda, \quad \left\| \frac{\partial \psi_i}{\partial x_2} \right\|_{L^2(\omega_0)}^2 \leq A_0(p) \lambda.$$

Multiplying the equation  $-\Delta \psi_i = \lambda_i \psi_i$  by  $\frac{\partial^2 \psi_i}{\partial x_1^2}$  and integrating by parts we find out that

$$\left\| \frac{\partial^2 \psi_i}{\partial x_1^2} \right\|_{L^2(\omega_1)}^2 \leq A_1(p) \lambda^2, \quad \left\| \frac{\partial^2 \psi_i}{\partial x_1 \partial x_2} \right\|_{L^2(\omega_1)}^2 \leq A_1(p) \lambda^2.$$

Hence (33) holds for  $n = 1$  and  $n = 2$ . Now assume that (33) holds for some  $n - 1$  and  $n$ . We will show that it holds for  $n + 1$  as well. Integration by parts yields

$$\begin{aligned} \left\| \Delta \left( \frac{\partial^{n-1} \psi_i}{\partial x_1^{n-1}} W_{n-1,p,\lambda} \right) \right\|_{L^2(\omega_{n-1})}^2 &= \left\| \frac{\partial^2}{\partial x_1^2} \left( \frac{\partial^{n-1} \psi_i}{\partial x_1^{n-1}} W_{n-1,p,\lambda} \right) \right\|_{L^2(\omega_{n-1})}^2 + \\ &\left\| \frac{\partial^2}{\partial x_2^2} \left( \frac{\partial^{n-1} \psi_i}{\partial x_1^{n-1}} W_{n-1,p,\lambda} \right) \right\|_{L^2(\omega_{n-1})}^2 + 2 \left\| \frac{\partial^2}{\partial x_1 \partial x_2} \left( \frac{\partial^{n-1} \psi_i}{\partial x_1^{n-1}} W_{n-1,p,\lambda} \right) \right\|_{L^2(\omega_{n-1})}^2 \end{aligned} \quad (34)$$

From the fact that  $W_{n-1,p,\lambda} = 1$  on the  $\omega_n$  it follows that the first and the last term on the right hand side of (34) are greater than or equal to

$$\left\| \frac{\partial^{n+1} \psi_i}{\partial x_1^{n+1}} \right\|_{L^2(\omega_n)}^2 \quad \text{and} \quad \left\| \frac{\partial^{n+1} \psi_i}{\partial x_1^n \partial x_2} \right\|_{L^2(\omega_n)}^2$$

respectively. The second term on the right hand side of (34) is positive and since  $\omega \subset \text{supp } W_{n,p,\lambda}$ , we get

$$\left\| \frac{\partial^{n+1} \psi_i}{\partial x_1^{n+1}} \right\|_{L^2(\omega_n)}^2 + \left\| \frac{\partial^{n+1} \psi_i}{\partial x_1^n \partial x_2} \right\|_{L^2(\omega_n)}^2 \leq \left\| \Delta \left( \frac{\partial^{n-1} \psi_i}{\partial x_1^{n-1}} W_{n-1,p,\lambda} \right) \right\|_{L^2(\omega_n)}^2. \quad (35)$$

Next we employ (32) and (33) to conclude that

$$\begin{aligned} \left\| \Delta \left( \frac{\partial^{n-1} \psi_i}{\partial x_1^{n-1}} W_{n-1,p,\lambda} \right) \right\|_{L^2(\omega_n)}^2 &= \\ \left\| \lambda_i \left( \frac{\partial^{n-1} \psi_i}{\partial x_1^{n-1}} \right) W_{n-1,p,\lambda} + \left( \frac{\partial^{n-1} \psi_i}{\partial x_1^{n-1}} \right) \Delta W_{n-1,p,\lambda} + 2 \left( \nabla \frac{\partial^{n-1} \psi_i}{\partial x_1^{n-1}} \right) \nabla W_{n-1,p,\lambda} \right\|_{L^2(\omega_n)}^2 & \\ \leq 3 \lambda^{n+1} A_{n-2}(p) + 6 \cdot 4^6 \alpha_2^2 p^4 \lambda^{n+1} A_{n-2}(p) + 24 \cdot 9^2 \alpha_1^2 p^2 \lambda^{n+1} A_{n-1}(p) &\leq \lambda^{n+1} A_n(p). \end{aligned} \quad (36)$$

□

As a consequence of this result we obtain

**Corollary 1.** *Let  $\omega$  be as in Lemma 1. Assume that  $\psi = \sum_{\lambda_i \leq \lambda} c_i \psi_i$  with  $\sum_{\lambda_i \leq \lambda} |c_i|^2 \leq V$ . Then*

$$\left\| \frac{\partial^{p+1} \psi}{\partial x_1^{p+1}} \right\|_{L^2(\omega)}^2 \leq \frac{A_p(p) V^2(\Omega)}{4\pi} \lambda^{p+2}.$$

*Proof.* By Lemma 1 and the Cauchy-Schwarz inequality we have

$$\left\| \frac{\partial^{p+1} \psi}{\partial x_1^{p+1}} \right\|_{L^2(\omega)}^2 \leq \sum_{\lambda_i \leq \lambda} |c_i|^2 \sum_{\lambda_i \leq \lambda} \left\| \frac{\partial^{p+1} \psi_i}{\partial x_1^{p+1}} \right\|_{L^2(\omega)}^2 \leq N_\lambda V A_p(p) \lambda^{p+1}, \quad (37)$$

Using the lower bound on  $\lambda_i$  given in (53) we find out that  $N_\lambda \leq \frac{V}{4\pi} \lambda$ .  $\square$

## 4.2 Lower bound on a square

Corollary 1 is one the two main technical results on which is based the proof of Theorems 1 and 2. The goal of this section is to prove the second one of these results, namely Proposition 2 (see page 15). We start with a couple of one dimensional estimates concerning smooth functions on an interval  $[0, l]$ . Unless otherwise stated,  $\|\cdot\|$  denotes the  $L^2$ -norm on  $[0, l]$ .

**Lemma 2.** *Let  $f \in C^{p+1}[0, l]$ ,  $p \in \mathbb{N}$ . Then*

$$\max |f^{(p)}|^2 \leq \frac{3}{2} \left( \frac{1}{l} \|f^{(p)}\|^2 + l \|f^{(p+1)}\|^2 \right).$$

*Proof.* Let  $\max |f^{(p)}| = |f^{(p)}(t_0)|$  with  $t_0 \in [0, l]$ . For any  $t \in [0, l]$  we have

$$f^{(p)}(t) = f^{(p)}(t_0) + \int_{t_0}^t f^{(p+1)}(\tau) d\tau.$$

Integrating with respect to  $t$  and using the Jensen inequality gives

$$\begin{aligned} l |f^{(p)}(t_0)|^2 &\leq \frac{3}{2} \int_0^l |f^{(p)}(t)|^2 dt + 3 \int_0^l \left( \int_{t_0}^t f^{(p+1)}(\tau) d\tau \right)^2 dt \\ &\leq \frac{3}{2} \|f^{(p)}\|^2 + 3 \int_0^l t \|f^{(p+1)}\|^2 dt = \frac{3}{2} \left( \|f^{(p)}\|^2 + l^2 \|f^{(p+1)}\|^2 \right). \end{aligned}$$

$\square$

**Lemma 3.** *Let  $f \in C^2 [0, \frac{1}{2} \lambda^{-1/2}]$  and real-valued. Then one of the following inequalities holds true:*

$$\max |f| \max |f''| \leq \frac{1}{4} \max |f'|^2 \quad (38)$$

$$\max |f'| \leq 32 \lambda^{\frac{1}{2}} \max |f| \quad (39)$$

*Proof.* Let  $m_i = \max |f^{(i)}|$ ,  $i \in \{0, 1, 2\}$  and let  $t_0 \in [0, \frac{1}{2} \lambda^{-1/2}]$  be such that  $m_1 = |f'(t_0)|$ . Without loss of generality we assume that  $t_0 < \frac{1}{4} \lambda^{-1/2}$ , otherwise we consider the interval  $[0, t_0]$  instead of  $[t_0, \frac{1}{2} \lambda^{-1/2}]$ . Assume that  $f'(t_0) = m_1$ . If

$$t_0 + \frac{m_1}{m_2} \leq \frac{1}{2} \lambda^{-1/2}, \quad (40)$$

then the Taylor theorem says that

$$m_0 \geq f \left( t_0 + \frac{m_1}{m_2} \right) \geq f(t_0) + m_1 \left( \frac{m_1}{m_2} \right) - \frac{m_2}{2} \left( \frac{m_1}{m_2} \right)^2 \geq -m_0 + \frac{m_1^2}{2m_2},$$

which implies (38). If, on the contrary,

$$t_0 + \frac{m_1}{m_2} > \frac{1}{2} \lambda^{-1/2}, \quad \text{then} \quad \frac{m_1}{m_2} > \frac{1}{2} \lambda^{-1/2} - t_0 > \frac{1}{4} \lambda^{-1/2}.$$

In this case we have

$$m_0 \geq f\left(t_0 + \frac{1}{8} \lambda^{-1/2}\right) \geq f(t_0) + m_1 \frac{1}{8} \lambda^{-1/2} - \frac{m_2}{128} \lambda^{-1},$$

which implies

$$m_1 \frac{1}{8} \lambda^{-1/2} - \frac{m_1}{32 \lambda^{-1/2}} \lambda^{-1} \leq 2m_0.$$

From here we conclude that

$$\frac{m_1}{m_0} \leq \frac{64}{3 \lambda^{-1/2}} \leq 32 \lambda^{\frac{1}{2}}.$$

The proof in the case  $f'(t_0) = -m_1$  is analogous.  $\square$

**Proposition 1.** *Let  $f \in C^p [0, \frac{1}{2} \lambda^{-1/2}]$ ,  $p \in \mathbb{N}$  and let  $f$  be real-valued. Then one of the following inequalities holds true:*

$$\max |f'| \leq 4^{p+\frac{1}{2}} \lambda^{\frac{1}{2}} \max |f| \quad (41)$$

$$\max |f'| \leq \left( \frac{\max |f^{(p)}|}{\max |f|} \right)^{\frac{1}{p}} 4^{p-\frac{1}{2}} \max |f|. \quad (42)$$

*Proof.* Let  $m_i = \max |f^{(i)}|$ ,  $i = 1, \dots, p$ . There are two possibilities. Either for all  $i \leq p$  holds

$$\frac{m_i}{m_{i-1}} \geq 32 \lambda^{\frac{1}{2}}, \quad (43)$$

or there exists  $i_0 \in [1, p]$ , such that

$$\forall i < i_0 \quad \frac{m_i}{m_{i-1}} \geq 32 \lambda^{\frac{1}{2}}, \quad \frac{m_{i_0}}{m_{i_0-1}} < 32 \lambda^{\frac{1}{2}}. \quad (44)$$

In the first case  $\frac{m_i}{m_{i-1}} > \frac{1}{4} \frac{m_{i-1}}{m_{i-2}}$  holds for all  $i \leq p$ , see Lemma 3. This yields

$$m_p \geq 4^{-\frac{p(p-1)}{2}} \left( \frac{m_1}{m_0} \right)^p m_0,$$

which is equivalent to (42). In the second case we have  $\frac{m_i}{m_{i-1}} > \frac{1}{4} \frac{m_{i-1}}{m_{i-2}}$  for all  $i \leq i_0$ . Combining this with (44) we conclude that

$$\frac{m_1}{m_0} \leq 4^{i_0+\frac{1}{2}} \lambda^{\frac{1}{2}}. \quad \square$$

**Corollary 2.** *Let  $f \in C^p [0, \frac{1}{2} \lambda^{-1/2}]$ ,  $p \in \mathbb{N}$  be a complex-valued function such that  $f(0) = 0$  and  $\max |f^{(p)}| \leq C(p) \lambda^{\frac{p}{2}+1}$  for some constant  $C(p)$ . Then for any  $\varphi_0, \varphi_1 \in \mathbb{R}$  holds*

$$\int_0^{\frac{1}{2} \lambda^{-1/2}} |f(t) - e^{i\varphi_1 t + i\varphi_0}|^2 dt \geq \frac{\lambda^{-\frac{1}{2}}}{9} \min \left\{ 4^{-p-\frac{5}{2}}, 4^{-\frac{p+3}{2}} 6^{\frac{1}{p}} C(p)^{-\frac{1}{p}} \lambda^{-\frac{1}{p}} \right\}. \quad (45)$$

*Proof.* Let  $u = \operatorname{Re} f$  and  $v = \operatorname{Im} f$ . If  $\max |f| \geq 6$ , then at least one of the expressions  $\max |u|$ ,  $\max |v|$  is larger than or equal to 3. Without loss of generality we assume that  $\max |u| \geq 3$  and apply Proposition 1 to the function  $u$ . If  $u$  satisfies (41), then there exists a subinterval  $I \subset [0, \frac{1}{2} \lambda^{-1/2}]$  of the length  $3^{-1} 4^{-p-\frac{1}{2}} \lambda^{-\frac{1}{2}}$  on which  $|u| \geq 3/2$ . This implies

$$\int_0^{\frac{1}{2} \lambda^{-1/2}} |f(t) - e^{i\varphi_1 t + i\varphi_0}|^2 dt \geq 3^{-1} 4^{-p-\frac{1}{2}} \lambda^{-\frac{1}{2}}.$$

If, on the other hand,  $u$  satisfies (42), then the length of the subinterval of  $[0, \frac{1}{2} \lambda^{-1/2}]$ , on which  $|u| \geq 3/2$ , is at least  $3^{-1} 4^{-\frac{p-1}{2}} C(p)^{-\frac{1}{p}} \lambda^{-\frac{1}{2}-\frac{1}{p}}$ , which gives

$$\int_0^{\frac{1}{2} \lambda^{-1/2}} |f(t) - e^{i\varphi_1 t + i\varphi_0}|^2 dt \geq 3^{-1} 4^{-\frac{p-1}{2}} C(p)^{-\frac{1}{p}} \lambda^{-\frac{1}{2}-\frac{1}{p}}.$$

Assume now that  $\max |f| < 6$ . The latter means that  $\max |u| < 6$  and  $\max |v| < 6$ . Since  $u(0) = v(0) = 0$ , there exists a subinterval of  $[0, \frac{1}{2} \lambda^{-1/2}]$ , on which  $\max\{|u(t)|, |v(t)|\} \leq 1/3$ , which implies  $|f(t) - e^{i\varphi_1 t + i\varphi_0}|^2 \geq 1/4$ . Applying Proposition 1 to the functions  $u, v$  we find out that the length of this interval is bounded from below by

$$\min \left\{ 3^{-2} 4^{-p-\frac{5}{2}} \lambda^{-\frac{1}{2}}, 3^{-2} 4^{-\frac{p+3}{2}} 6^{\frac{1}{p}} C(p)^{-\frac{1}{p}} \lambda^{-\frac{1}{2}-\frac{1}{p}} \right\}.$$

This completes the proof.  $\square$

With the above auxiliary results at hand, we can finally prove the following integral estimate, which will play a central role in the proof of Theorem 1 and 2.

**Proposition 2.** *Let  $f \in C^{p+1}[\omega]$  a complex valued function such that  $f(0, x_2) = 0$  for each  $x_2$  and*

$$\left\| \frac{\partial^{p+1} f}{\partial x_1^{p+1}} \right\|_{L^2(\omega)} \leq \beta_{p+1} \lambda^{1+\frac{p}{2}}, \quad \left\| \frac{\partial^p f}{\partial x_1^p} \right\|_{L^2(\omega)} \leq \beta_p \lambda^{\frac{1}{2}+\frac{p}{2}}$$

for some positive  $\beta_p$  and  $\beta_{p+1}$ . Then the inequality

$$\left\| f - e^{i(\xi_1 x_1 + \xi_2 x_2 + \varphi)} \right\|_{L^2(\omega)}^2 \geq \frac{1}{36} \min \left\{ 4^{-p-\frac{5}{2}} \lambda^{-1}, 4^{-\frac{p}{2}-\frac{3}{2}} 6^{\frac{1}{2p}} (\beta_{p+1}^2 + \beta_p^2)^{-\frac{1}{2p}} \lambda^{-1-\frac{1}{p}} \right\} \quad (46)$$

holds true for all  $\xi_1, \xi_2, \varphi \in \mathbb{R}$ .

*Proof.* The measure of the set

$$\left\{ x_2 \in [0, \lambda^{-1/2}] : \int_0^{\frac{1}{2\sqrt{x}}} \left| \frac{\partial^i f(x_1, x_2)}{\partial x_1^i} \right|^2 dx_1 \leq 8 \beta_i^2 \lambda^{i+\frac{3}{2}}, i \in \{p, p+1\} \right\}$$

is obviously at least  $\frac{1}{4} \lambda^{-\frac{1}{2}}$ . For such  $x_2$  holds by Lemma 2

$$\max_{x_1} \left| \frac{\partial^p f(x_1, x_2)}{\partial x_1^p} \right| \leq \sqrt{3} \lambda^{1+\frac{p}{2}} \sqrt{\beta_{p+1}^2 + \beta_p^2}.$$

Corollary 2 then implies the statement.  $\square$

### 4.3 Proof of Theorem 1

*Proof of Theorem 1.* Fix  $\lambda > 0$ . Let  $\lambda_j$  be the eigenvalues of the Dirichlet Laplacian on  $\Omega$  and let  $\psi_j$  be the corresponding normalised eigenfunctions. For  $k \in \mathbb{N}$  we define

$$F(\xi) = \sum_{j=1}^k |\hat{\psi}_j(\xi)|^2,$$

where  $\hat{\psi}_j$  denotes the Fourier transform of  $\psi_j$ . Moreover, we denote by  $F^*(|\xi|)$  the decreasing radial rearrangement of  $F(\xi)$ . Let

$$\psi(x) = \sum_{\lambda_i \leq \lambda} c_i \psi_i(x), \quad \text{with} \quad \sum_{\mu_i \leq \lambda} |c_i|^2 \leq V.$$

For each  $j = 1, \dots, n$  we choose on the middle part of  $p_j$  several points  $t_l$  such that  $\text{dist}(t_l, t_{l+1}) = \sqrt{2} \lambda^{-1/2}$  for all  $l$  and denote by  $T_l$  the squares with the side  $\frac{1}{2} \lambda^{-1/2}$  constructed in the middle point between  $t_l$  and  $t_{l+1}$ , see Figure 2. We note that for each  $j$  the number of these squares is at least

$$N_j = \left\lceil \frac{1}{3\sqrt{2}} l_j \lambda^{\frac{1}{2}} \right\rceil.$$

According to Corollary 1 for each  $l$  and  $p$  we have

$$\left\| \frac{\partial^{p+1} \psi}{\partial \nu^{p+1}} \right\|_{L^2(T_l)}^2 \leq \frac{A_p(p) V^2}{4\pi} \lambda^{p+2},$$

where  $\frac{\partial \psi}{\partial \nu}$  denotes the normal derivative of  $\psi$ . In view of Proposition 2 and Corollary 3 we get

$$\|\psi - e^{i\xi \cdot x}\|_{L^2(T_l)}^2 \geq \frac{1}{36} \min \left\{ 4^{-p-\frac{5}{2}} \lambda^{-1}, 4^{-\frac{p}{2}-\frac{3}{2}} 6^{\frac{1}{2p}} (\beta_{p+1}^2 + \beta_p^2)^{-\frac{1}{2p}} \lambda^{-1-\frac{1}{p}} \right\}, \quad (47)$$

where

$$\beta_{p+1}^2 = \frac{A_p(p) V^2}{4\pi}.$$

We continue by estimating the sequence  $A_p(p)$ . A direct inspection shows that

$$A_p(p) \leq c_0 2^{(p+1)^2}, \quad c_0 = 7 \cdot 10^{22}. \quad (48)$$

This implies that  $(\beta_{p+1}^2 + \beta_p^2)^{-\frac{1}{2}} \geq 2^{-1/2} \sqrt{\pi} c_0^{-1/2} V^{-1} 2^{-(p+1)^2/2}$ . Hence for

$$p = \left\lceil \sqrt{2 \log_2(V\lambda/c_1)} \right\rceil - 1, \quad c_1 = \sqrt{\frac{3\pi}{2}} c_0^{-\frac{1}{2}}$$

we obtain

$$\|\psi - e^{i\xi \cdot x}\|_{L^2(T_l)}^2 \geq \frac{2^{-3}}{36} c_1^{-1} V \left( \frac{V\lambda}{c_1} \right)^{-1 - \frac{2}{\sqrt{\log_2(V\lambda/c_1)}}}.$$

Taking  $\lambda$  large enough such that

$$\lambda^{-1/2} \leq \frac{d_j}{3}.$$

we make sure that the squares  $T_l$  lie inside  $\Omega$  and that they do not overlap each other. Summing this inequality for all  $l = 1, \dots, N_j$  and all  $j = 1, \dots, n$  we thus arrive at

$$V - 4\pi^2 F^*(|\xi|) \geq \|\psi - e^{i\xi \cdot x}\|_{L^2(\Omega^c)}^2 \geq c_2 V^{\frac{1}{2}} \left( \frac{V\lambda}{c_1} \right)^{-\frac{1}{2} - \frac{2}{\sqrt{\log_2(V\lambda/c_1)}}} \sum_{j=1}^n l_j \Theta \left( \lambda - \frac{9}{d_j^2} \right) \quad (49)$$

with  $c_2 = \frac{2^{-3}}{9\sqrt{2}36} c_1^{-1/2}$ . This yields the following upper bound on  $F^*$ :

$$F^*(|\xi|) \leq M(p, \lambda) := \frac{V}{4\pi^2} \left[ 1 - c_2 V^{-\frac{1}{2}} \left( \frac{V\lambda}{c_1} \right)^{-\frac{1}{2} - \frac{2}{\sqrt{\log_2(V\lambda/c_1)}}} \sum_{j=1}^n l_j \Theta \left( \lambda - \frac{9}{d_j^2} \right) \right]. \quad (50)$$

Now we use the minimiser (10) with  $V/4\pi^2$  replaced by  $M(p, \lambda)$  to obtain

$$\sum_{j=1}^k \lambda_j \geq \int_{\mathbb{R}^2} F^*(|\xi|) |\xi|^2 d\xi \geq \frac{\lambda^2 V^2}{8\pi^3 M(p, \lambda)}. \quad (51)$$

Employing the definition of  $M(p, \lambda)$  we then find out that

$$\sum_{j=1}^k \lambda_j \geq \frac{\lambda^2 V}{2\pi} + c_2 c_1^2 V^{-\frac{3}{2}} \left( \frac{V\lambda}{c_1} \right)^{\frac{3}{2} - \frac{2}{\sqrt{\log_2(V\lambda/c_1)}}} \sum_{j=1}^n l_j \Theta \left( \lambda - \frac{9}{d_j^2} \right). \quad (52)$$

Next we set  $\lambda = \lambda_k$  and note that inequality (2) yields

$$\frac{2\pi}{V} k \leq \lambda_k. \quad (53)$$

Since the right hand side of (52) is an increasing function of  $\lambda$ , we can use (53) to conclude that

$$\sum_{j=1}^k \lambda_j \geq \frac{2\pi}{V} k^2 + 4c_3 k^{\frac{3}{2} - \frac{2}{\sqrt{\log_2(2\pi k/c_1)}}} \sum_{j=1}^n l_j \Theta \left( k - \frac{9V}{2\pi d_j^2} \right) V^{-3/2} \quad (54)$$

where

$$c_3 = \frac{2^{-3}}{9\sqrt{2}36} (2\pi)^{\frac{5}{4}} c_1^{1/4}.$$

Finally, we combine inequalities (54) and (16) to get (21).  $\square$

## 5 Proof for general domains

From now on we suppose that  $\Omega$  is a general domain satisfying assumption A. To prove a Li-Yau type inequality with the correction term we cannot directly employ the approach invented for polygons, since  $\partial\Omega$  is in general nowhere straight. However, we can extend  $\Omega$  by adding small ‘‘bumps’’ to certain parts of  $\partial\Omega$ , see Figure 4, in order to obtain an extended domain  $\Omega^e$  whose boundary is in certain parts represented by a straight line. On these straight pieces of  $\partial\Omega^e$  we will then employ the same strategy as in the case of polygons. Due to the monotonicity of eigenvalues, any lower bound on the sum of the eigenvalues on the extended domain gives also a lower bound on the sum of the eigenvalues on  $\Omega$ . On the other hand, we have to make sure that the volume of  $\Omega^e$  is not much bigger than  $V$ , because otherwise it could destroy the effect of the correction term in (26) by decreasing the leading term. We will again split the exposition in several steps.

### 5.1 Step 1: Some geometrical remarks

Here we will show that  $\partial\Omega \cap \Gamma_j$  can be locally represented as a graph of a certain  $C^2$ -smooth function. Let  $\Gamma = \{x_1(s), x_2(s)\}$  be a part of the boundary of  $\Omega$  parametrised by its length  $s$  and such that  $x_1(s), x_2(s) \in C^2(\mathbb{R}_+)$ . Let

$$\varkappa_0 := \max_{\{x_1, x_2\} \in \Gamma} |\varkappa(x_1, x_2)|$$

be the maximal curvature of  $\Gamma$ . We consider certain points  $A = \{x_1(s'), x_2(s')\} \in \Gamma$  and  $B = \{x_1(s''), x_2(s'')\} \in \Gamma$  and chose a new system  $(u, v)$  such that  $A = (0, 0)$  and the  $u$ -axes goes along the line  $AB$ .

**Lemma 4.** *Assume that  $\varkappa_0 |s' - s''| \leq \pi/4$ . Then the following statements hold true.*

(i) *The part of  $\Gamma$  connecting  $A$  and  $B$  can be written in the system of coordinates  $(u, v)$  as  $v = v(u)$ ,  $u \in [0, u_0]$ , where  $u_0 = |AB|$ . Moreover, we have*

$$\max_{u \in [0, u_0]} v(u) \leq \sqrt{2} \varkappa_0 u_0^2. \quad (55)$$

(ii) *The inequality*

$$2^{-1/2} |s' - s''| \leq |AB| \leq |s' - s''| \quad (56)$$

*holds.*

*Proof.* Let  $\{u(s), v(s)\}$  be the parametrisation of  $\Gamma$  in the coordinates  $(u, v)$ . By assumption we have

$$\int_0^{|s' - s''|} \varkappa(s) ds \leq \pi/4 \quad (57)$$

This means that for any  $s \in [0, |s' - s''|]$  the angle between the tangent of  $\Gamma$  at the point  $\{u(s), v(s)\}$  and the  $u$ -axes is less than or equal to  $\pi/4$ . Assume that there exists  $s_1, s_2 \in [0, |s' - s''|]$  such that  $u(s_1) = u(s_2)$ . Then there exists  $s_3 \in [s_1, s_2]$  such that the tangent of  $\Gamma$  at  $\{u(s_3), v(s_3)\}$  is orthogonal to the  $u$ -axes. The latter contradicts (57). This shows that the part of  $\Gamma$  between  $A$  and  $B$  can be considered as the graph of the function

$$v = v(u), \quad u \in [0, u_0], \quad v(0) = v(u_0) = 0.$$

This proves the first part of (i) and, in view of (57), shows that  $|v'(u)| \leq 1$  on  $[0, u_0]$ . Next we prove inequality (56). It thus follows that

$$u_0 = |AB| \leq |s' - s''| = \int_0^{u_0} (1 + |v'(u)|^2)^{1/2} du \leq 2^{1/2} u_0,$$

which implies (56). To prove (55) we note that  $v(u)$  is twice differentiable and therefore there exists some  $u_1 \in [0, u_0]$ , such that  $v'(u_1) = 0$ . Since  $|v''(u)| = |\varkappa(u)| (1 + |v'(u)|^2)^{3/2} \leq 2^{3/2} \varkappa_0$ , we obtain

$$|v'(u)| \leq \int_{u_1}^u |v''(u)| du \leq 2^{3/2} \varkappa_0 u_0 \quad \forall u \in [0, u_0].$$

The last inequality together with the fact that  $v(0) = v(u_0) = 0$  finally implies

$$|v(u)| \leq \frac{1}{2} 2^{3/2} \varkappa_0 u_0^2 = 2^{1/2} \varkappa_0 u_0^2 \quad \forall u \in [0, u_0].$$

□

## 5.2 Step 2: Approximation of the boundary

Next we introduce a procedure that allows us to choose appropriate parts of  $\partial\Omega \cap \Gamma_j$  on which we will construct the additional “bumps”, see Figure 4. Let  $\Gamma_j$ ,  $j = 1 \dots m$  be the parts of boundary defined

in section 3 with the end points  $A_j, B_j$  and the partition  $a_i^j, i = 0, \dots, n_j$ . We fix  $j \in \{1, \dots, m\}$  and take  $\lambda$  large enough, such that

$$\lambda^{-\frac{1}{2}} \leq \min \left\{ \frac{d_j}{3}, \frac{\pi}{8\sqrt{2}\varkappa_j} \right\}, \quad \text{if } L(\Gamma_j) > \frac{3\pi}{8\varkappa_j} \quad (58)$$

and

$$\lambda^{-\frac{1}{2}} \leq \min \left\{ \frac{d_j}{3}, \frac{L(\Gamma_j)}{3\sqrt{2}} \right\}, \quad \text{if } L(\Gamma_j) \leq \frac{3\pi}{8\varkappa_j}. \quad (59)$$

Let us consider  $\Gamma_j \cap (a_i^j, a_{i+1}^j)$  with  $0 < i < n_j$ . On this part of the boundary we choose several disjoint arcs  $(b_l, b'_l)$ , such that each of them has the length  $\sqrt{2}\lambda^{-1/2}$  and such that

$$\sum_l s(b_l, b'_l) \geq \frac{1}{3} s(a_i^j, a_{i+1}^j), \quad s(a_i^j, a_{i+1}^j) - \sum_l s(b_l, b'_l) \leq \sqrt{2}\lambda^{-1/2},$$

where  $s(a, b)$  denotes the arc-length between  $a$  and  $b$ .

Next we pick an  $l$  and connect  $b_l$  and  $b'_l$  with a straight line and choose a local system of coordinates  $(y_1, y_2)$  so that the  $y_1$ -axis goes along the straight line from  $b_l$  to  $b'_l$  and the origin is in  $b_l$ . Notice that  $s(a_{i-1}^j, a_{i+1}^j) = s(a_i^j, a_{i+2}^j) \leq \frac{\pi}{2\varkappa_j}$ , which according to Lemma 4 means that in the chosen coordinate system the boundary between  $a_{i-1}^j$  and  $a_{i+2}^j$  can be written explicitly as  $y_2 = f(y_1)$ . Let  $y_0 = \text{dist}(b_l, b'_l)$ . In view of Lemma 4

$$\max_{y_1} |f(y_1)| \leq \sqrt{2}\varkappa_j y_0^2 \leq 2^{\frac{3}{2}}\varkappa_j \lambda^{-1}.$$

Now we introduce

$$\Sigma_1 = \left\{ (y_1, y_2) : y_1 \in [0, y_0], y_2 = 2^{\frac{3}{2}}\varkappa_j \lambda^{-1} \right\}$$

and

$$\Sigma_2 = \left\{ (y_1, y_2) : y_1 \in [0, y_0], y_2 = -2^{\frac{3}{2}}\varkappa_j \lambda^{-1} \right\}$$

**Lemma 5.** *If  $\lambda > 6\varkappa_j/d_j$ , then*

$$\Sigma_1 \cap \partial\Omega = \Sigma_2 \cap \partial\Omega = \emptyset.$$

*Proof.* Obviously  $\Sigma_1$  and  $\Sigma_2$  do not cross  $\partial\Omega$  between  $a_{i-1}^j$  and  $a_{i+2}^j$ . On the other hand, for each point  $P = (y_1^P, y_2^P)$  holds

$$\text{dist}(P, (a_i^j, a_{i+1}^j)) \leq 2^{3/2}\varkappa_j \lambda^{-1}.$$

Since  $\text{dist}\left((a_i^j, a_{i+1}^j), \partial\Omega \setminus (a_{i-1}^j, a_{i+2}^j)\right) \geq d_j$ , this implies

$$\text{dist}\left(P, \partial\Omega \setminus (a_{i-1}^j, a_{i+2}^j)\right) \geq d_j - 2^{3/2}\varkappa_j \lambda^{-1} > \frac{d_j}{2} > 0.$$

□

The last Lemma says that one of the sets  $\Sigma_1$  and  $\Sigma_2$  is inside  $\Omega$  and the other one is outside  $\Omega$ . Without loss of generality we assume that  $\Sigma_1$  is outside  $\Omega$ .

### 5.3 Step 3: Extended domain $\Omega^e$ .

The extended domain  $\Omega^e$  differs from  $\Omega$  if  $\lambda$  is large enough so that (58) respectively (59) is satisfied (otherwise it coincides with  $\Omega$ ).

To define  $\Omega^e$  we proceed as follows. For a fixed  $j \in \{1, \dots, m\}$ , fixed  $i \in \{1, \dots, n_j - 1\}$  and fixed  $l$ , we consider the boundary between the points  $b_l$  and  $b'_l$ . If it is a straight line, we do not change it. Otherwise we replace this piece of the boundary with the segment  $\Sigma_i$ , where  $i$  is such that  $\Sigma_i$  is outside  $\Omega$ , and connect the end points of  $\Sigma_1$  with the boundary at certain points  $\tilde{b}_l \in (b'_{l-1}, b_l)$  and  $\tilde{b}'_l \in (b'_l, b_{l+1})$  with appropriate  $C^2$  functions. We choose these function and the points  $\tilde{b}_l, \tilde{b}'_l$  in such a way that the added area to  $\Omega$  is less than 3 times the area of the rectangle with the corners given by  $b_l, b'_l$  and the end points of  $\Sigma_1$ . We then obtain a new region whose boundary, corresponding to the original piece  $\Gamma_j$  is again  $C^2$ -smooth and which between the original boundary points  $b_l$  and  $b'_l$  consists of a straight line. Repeating this procedure for all  $\Gamma_j, j = 1, \dots, m$ , all  $i \in \{1, \dots, n_j - 1\}$  and all  $l$  we thus obtain a new domain  $\Omega^e$ .

As a next step we construct the squares  $T_l$  of the side  $\frac{1}{2}\lambda^{-1/2}$  between the the points  $b_l$  and  $b'_l$  centred in the middle. Note that, according to Lemma 4,  $|b_l b'_l| \geq \lambda^{-1/2}/\sqrt{2}$ . We have

**Lemma 6.** *The squares  $T_l$  do not overlap.*

*Proof.* First we show that every  $T_l$  does not overlap with any of the squares constructed on the part of the boundary different from the arch  $(a_{i-1}^j, a_{i+2}^j)$ . Indeed, each point of  $T_l$  has distance to  $(b_l, b'_l)$  at most  $\frac{1}{2}\lambda^{-1/2}$  and the distance between  $(b_l, b'_l)$  and  $\partial\Omega \setminus (a_{i-1}^j, a_{i+2}^j)$  is at least  $d_j$ . Since  $\lambda^{-1/2} < d_j$ , see (58), the result follows.

Consider now  $(a_{i-1}^j, a_{i+2}^j)$ . This part can be written as  $y_2 = f(y_1)$  in the above introduced coordinate system. Consider the squares  $T_{l_1}$  and  $T_{l_2}$  with  $l_1 \neq l_2$ . Let  $y_1^1$  be the  $y_1$  coordinate of the middle point between  $b_{l_1}$  and  $b'_{l_1}$  and let  $y_1^2$  be the  $y_1$  coordinate of the middle point between  $b_{l_2}$  and  $b'_{l_2}$ . Since  $|f'(y_1)| \leq 1$  on  $(a_{i-1}^j, a_{i+2}^j)$ , we have  $|y_1^1 - y_1^2| \geq \lambda^{-1/2}$ . For all points  $(y_1, y_2) \in T_{l_1}$  holds  $|y_1 - y_1^1| \leq \frac{1}{4}\lambda^{-1/2}$  and for all points  $(y_1, y_2) \in T_{l_2}$  holds  $|y_1 - y_1^2| \leq \frac{\sqrt{2}}{2}\lambda^{-1/2}$ . Collecting these inequalities we conclude that  $T_{l_1} \cap T_{l_2} = \emptyset$ .  $\square$

As a consequence of the last result we obtain estimates on the volume of  $\Omega^e$ , which will be used in the proof of Theorem 2.

**Corollary 3.** *Let  $V^e$  be the volume of the extended domain  $\Omega^e$ . Then*

$$V^e \leq V + 2^{\frac{3}{2}} \lambda^{-1} \sum_{j=1}^m \varkappa_j L(\Gamma_j). \quad (60)$$

Moreover, if

$$\lambda \geq \Lambda_1 := 9 \cdot 2^{10} \max_j \varkappa_j^2,$$

then

$$V^e \leq 2V. \quad (61)$$

*Proof.* Inequality (60) follows directly from the construction of  $\Omega^e$ , since the area of the added volume along  $\Gamma_j$  does not exceed  $2^{\frac{3}{2}} \lambda^{-1} \varkappa_j L(\Gamma_j)$ . As for the second inequality, we consider each pair  $b_l, b'_l$  and note that for  $\lambda \geq 9 \cdot 2^{10} \varkappa_j^2$  is the area of the added volume between  $\tilde{b}_l$  and  $\tilde{b}'_l$ , bounded from above by

$$12 \varkappa_j \lambda^{-\frac{3}{2}} \leq \frac{1}{8} \lambda^{-1}.$$

This follows from the choice of the points  $b_l$ , see section 5.3. On the other hand, for  $\lambda$  chosen as above we get

$$|T_l \cap \Omega| \geq \frac{1}{2} |T_l| = \frac{1}{8} \lambda^{-1}.$$

Since  $T_l$  do not overlap, we obtain (61).  $\square$

## 5.4 Proof of Theorem 2

*Proof of Theorem 2.* Fix  $\lambda > 0$  and consider the extended domain  $\Omega^e$ . Let  $\mu_j$  be the eigenvalues of the Dirichlet Laplacian on  $\Omega^e$  and let  $\phi_j$  be the corresponding normalised eigenfunctions. For  $k \in \mathbb{N}$  we define

$$F_e(\xi) = \sum_{j=1}^k |\hat{\phi}_j(\xi)|^2,$$

where  $\hat{\phi}_j$  denotes the Fourier transform of  $\phi_j$ . By  $F_e^*(|\xi|)$  we denote the decreasing radial rearrangement of  $F_e(\xi)$ . Let

$$\phi(x) = \sum_{\mu_i \leq \lambda} c_i \phi_i(x), \quad \text{with} \quad \sum_{\mu_i \leq \lambda} |c_i|^2 \leq V^e$$

and let  $T_l$  be the sequence of squares constructed along  $\Gamma_j$ . For each  $j$  is the number of these squares at least

$$N_j = \left\lceil \frac{1}{9\sqrt{2}} L(\Gamma_j) \lambda^{\frac{1}{2}} \right\rceil.$$

Next we take  $\lambda \geq \Lambda_1$ , so that  $V^e \leq 2V$ , see Corollary 3. According to Corollary 1 for each  $l$  and  $p$  we then have

$$\left\| \frac{\partial^{p+1} \phi}{\partial \nu^{p+1}} \right\|_{L^2(R_n)}^2 \leq \frac{A_p(p)(V^e)^2}{4\pi} \lambda^{p+2} \leq \frac{A_p(p)V^2}{\pi} \lambda^{p+2},$$

where  $\frac{\partial \phi}{\partial \nu}$  denotes the normal derivative of  $\phi$ . In view of Proposition 2 for each  $l$  holds

$$\|\phi - e^{i\xi \cdot x}\|_{L^2(T_l)}^2 \geq \frac{1}{36} \min \left\{ 4^{-p-\frac{5}{2}} \lambda^{-1}, 4^{-\frac{p}{2}-\frac{3}{2}} 6^{\frac{1}{2p}} (\beta_{p+1}^2 + \beta_p^2)^{-\frac{1}{2p}} \lambda^{-1-\frac{1}{p}} \right\},$$

with

$$\beta_{p+1}^2 = \frac{A_p(p)(V^e)^2}{4\pi} \leq \frac{A_p(p)V^2}{\pi}.$$

Now we employ the same arguments used in the proof of Theorem 1 in order to find an appropriate upper bound on  $F_e^*$ . Since  $\lambda \geq \Lambda_1$  we can use Corollary 3 to arrive at

$$F_e^*(|\xi|) \leq \frac{V}{4\pi^2} \left[ 1 + \sum_{j=1}^m \left( 2^{3/2} V^{-1} \varkappa_j \lambda^{-1} - \frac{c_2}{2} V^{-\frac{1}{2}} \left( \frac{V\lambda}{c_1} \right)^{-\frac{1}{2} - \frac{2}{\sqrt{\log_2(V\lambda/c_1)}}} \right) L(\Gamma_j) \right].$$

Note that for

$$\lambda \geq \Lambda_2 := 2^{2^6} c_1 V^{-1}$$

we have  $\left( \frac{V\lambda}{c_1} \right)^{-\frac{1}{2} - \frac{2}{\sqrt{\log_2(V\lambda/c_1)}}} \geq \left( \frac{V\lambda}{c_1} \right)^{-\frac{3}{4}}$  and therefore

$$F_e^*(|\xi|) \leq M_e(p, \lambda) := \frac{V}{4\pi^2} \left[ 1 - \frac{c_2}{4} V^{-\frac{1}{2}} \left( \frac{V\lambda}{c_1} \right)^{-\frac{1}{2} - \frac{2}{\sqrt{\log_2(V\lambda/c_1)}}} \sum_{j=1}^m L(\Gamma_j) \Theta(\lambda - \Lambda_3(j)) \right].$$

where

$$\Lambda_3(j) := \max \{ \Lambda_1, \Lambda_2, c_1^{-1} 2^{2^2} 6^8 \varkappa_j^4 V \}.$$

We now use again the Li-Yau type minimiser (10) with  $V/4\pi^2$  replaced by  $M_e(p, \lambda)$  to obtain

$$\sum_{j=1}^k \lambda_j \geq \sum_{j=1}^k \mu_j \geq \int_{\mathbb{R}^2} F_e^*(|\xi|)|\xi|^2 d\xi \geq \frac{\lambda^2 V^2}{8\pi^3 M_e(p, \lambda)}.$$

As in the proof of Theorem 1 we set  $\lambda = \lambda_k$  and use definition of  $M_e(p, \lambda)$  together with inequalities (58), (59) and (53) to obtain

$$\sum_{j=1}^k \lambda_j \geq \frac{2\pi}{V} k^2 + c_3 k^{\frac{3}{2} - \frac{2}{\sqrt{\log_2(2\pi k/c_1)}}} \sum_{j=1}^m L(\Gamma_j) \Theta(k - k(j)) V^{-3/2} \quad (62)$$

where

$$k(j) := \frac{V}{2\pi} \max \left\{ \Lambda_3(j), \frac{9}{d_j^2}, \frac{128 \kappa_j^2}{\pi^2}, \frac{6\epsilon_j}{d_j} \right\}.$$

Finally, we combine inequalities (62) and (16) to get (26).  $\square$

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